

Stochastic Partial Differential Equations in Fluid Mechanics

Lecture 2: The Navier-Stokes equations with rough force

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The Navier-Stokes equations

Assume D is a regular bounded connected open domain. Consider the equations with *non-differentiable* forcing term W :

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + f + \partial_t W \\ \operatorname{div} u &= 0\end{aligned}$$

supplemented by boundary and initial condition

$$\begin{aligned}u|_{\partial D} &= 0 \\ u|_{t=0} &= u_0.\end{aligned}$$

$$\begin{aligned}H &= \{v \in L^2(D; \mathbb{R}^2) : \operatorname{div} v = 0, v \cdot n|_{\partial D} = 0\} \\V &= \{v \in H^1(D; \mathbb{R}^2) : \operatorname{div} v = 0, v|_{\partial D} = 0\} \\D(A) &= \{v \in H^2(D; \mathbb{R}^2) : \operatorname{div} v = 0, v|_{\partial D} = 0\} \\D(A) &\subset V \subset H \cong H' \subset V'\end{aligned}$$

$P : L^2(D; \mathbb{R}^2) \rightarrow H$ projection

$$w = v + \nabla q \quad w \in L^2(D; \mathbb{R}^2), v = Pw \in H, \langle v, \nabla q \rangle = 0$$

$$A : D(A) \rightarrow H, \quad Av = vP\Delta v$$

$$\mathbb{L}^4 = L^4(D, \mathbb{R}^2) \cap H$$

$$b : \mathbb{L}^4 \times V \times \mathbb{L}^4 \rightarrow \mathbb{R}$$

$$b(u, v, w) = \sum_{i,j=1}^2 \int_D u_i(x) \partial_i v_j(x) w_j(x) dx = \int_D (u \cdot \nabla v) \cdot w dx$$

$$B : \mathbb{L}^4 \times \mathbb{L}^4 \rightarrow V'$$

$$\langle B(u, v), \phi \rangle = -b(u, \phi, v) = - \int_D (u \cdot \nabla \phi) \cdot v dx$$

$$B(u, v) = P(u \cdot \nabla v)$$

Definition

Given $u_0 \in H$, $f \in L^2(0, T; V')$ and $W \in L^\infty(0, T; D(A))$, we say that

$$u \in C([0, T]; H) \cap L^\infty(0, T; \mathbb{L}^4) \\ + C([0, T]; H) \cap L^2(0, T; V)$$

is a weak solution of the NS equations if

$$u - z \in C([0, T]; H) \cap L^2(0, T; V)$$

where z is defined above with $z_0 = 0$ and

$$\langle u(t), \phi \rangle - \int_0^t b(u(s), \phi, u(s)) ds \\ = \langle u_0, \phi \rangle + \int_0^t \langle u, A\phi \rangle ds + \int_0^t \langle f, \phi \rangle ds + \langle W(t), \phi \rangle - \langle W(0), \phi \rangle$$

for every $\phi \in D(A)$.

Theorem

Assume $u_0 \in H$, $f \in L^2(0, T; V')$ and $W \in L^\infty(0, T; D(A))$. Then the Navier-Stokes equations have a unique weak solution, given by the sum of the solution.

The solution of the Navier-Stokes equations is given by the sum of the solution of Stokes problem and the solution of the auxiliary problem.

Stokes problem

Definition

Given $z_0 \in H$ and $W \in L^\infty(0, T; D(A))$, we say that z is a weak solution of Stokes problem if

$$z \in L^\infty(0, T; H)$$

and

$$\langle z(t), \phi \rangle = \langle z_0, \phi \rangle + \int_0^t \langle z(s), A\phi \rangle ds + \langle W(t), \phi \rangle - \langle W(0), \phi \rangle$$

for every $\phi \in D(A)$.

Theorem

If $z_0 \in H$ and $W \in L^\infty(0, T; D(A))$, then there exists one and only one weak solution of Stokes problem; it is given by

$$z(t) = e^{tA} z_0 + W(t) - e^{tA} W(0) + \int_0^t e^{(t-s)A} A W(s) ds.$$

Step 1 (uniqueness and explicit formula). Let z be a solution. Let

$$\phi \in C^1([0, T]; H) \cap C([0, T]; D(A))$$

be given. Then

$$\begin{aligned} & \langle z(t), \phi(t) \rangle - \langle z_0, \phi(0) \rangle \\ = & \int_0^t \langle z(s), \partial_s \phi(s) \rangle ds + \int_0^t \langle z(s), A\phi(s) \rangle ds \\ & + \langle W(t), \phi(t) \rangle - \langle W(0), \phi(0) \rangle - \int_0^t \langle W(s), \partial_s \phi(s) \rangle ds. \end{aligned}$$

Take the function

$$\phi_t(s) := e^{(t-s)A}\psi$$

with $\psi \in D(A^2)$.

$$\begin{aligned}
& \langle z(t), \phi(t) \rangle - \langle z_0, \phi(0) \rangle \\
= & \int_0^t \langle z(s), \partial_s \phi(s) \rangle ds + \int_0^t \langle z(s), A\phi(s) \rangle ds \\
& + \langle W(t), \phi(t) \rangle - \langle W(0), \phi(0) \rangle - \int_0^t \langle W(s), \partial_s \phi(s) \rangle ds.
\end{aligned}$$

$$\phi_t(s) := e^{(t-s)A}\psi$$

$$\begin{aligned}
& \langle z(t), \psi \rangle - \langle z_0, e^{tA}\psi \rangle \\
= & - \int_0^t \langle z(s), Ae^{(t-s)A}\psi \rangle ds + \int_0^t \langle z(s), Ae^{(t-s)A}\psi \rangle ds \\
& + \langle W(t), \psi \rangle - \langle W(0), e^{tA}\psi \rangle + \int_0^t \langle W(s), Ae^{(t-s)A}\psi \rangle ds.
\end{aligned}$$

$$z(t) = e^{tA}z_0 + W(t) - e^{tA}W(0) + \int_0^t e^{(t-s)A}AW(s) ds.$$

\Rightarrow uniqueness and explicit formula.

Step 2 (existence).

$$z(t) = e^{tA} z_0 + W(t) - e^{tA} W(0) + \int_0^t e^{(t-s)A} A W(s) ds.$$

$$z(t) - W(t) = e^{tA} (z_0 - W(0)) + \int_0^t e^{(t-s)A} A W(s) ds$$

$$\frac{d}{dt} (z(t) - W(t)) = Az(t) - W(t) + AW(t)$$

$$\langle z(t), \phi \rangle = \langle z_0, \phi \rangle + \int_0^t \langle z(s), A\phi \rangle ds + \langle W(t), \phi \rangle - \langle W(0), \phi \rangle.$$

$$z(t) = e^{tA} z_0 + W(t) - e^{tA} W(0) + \int_0^t A e^{(t-s)A} W(s) ds.$$

Theorem

If $z_0 \in D\left((-A)^{\frac{1}{4}+\epsilon}\right)$ and

$$W \in L^\infty\left(0, T; D\left((-A)^{\frac{1}{4}+\epsilon}\right)\right)$$

for some $\epsilon > 0$ then

$$z \in L^\infty\left(0, T; D\left((-A)^{\frac{1}{4}+\frac{\epsilon}{2}}\right)\right) \subset L^\infty(0, T; \mathbb{L}^4).$$

$$z_0 \in D\left((-A)^{\frac{1}{4}+\epsilon}\right)$$

$$W \in L^\infty\left(0, T; D\left((-A)^{\frac{1}{4}+\epsilon}\right)\right)$$

$$z(t) = \underbrace{e^{tA}z_0 + W(t) - e^{tA}W(0)}_{\in L^\infty\left(0, T; D\left((-A)^{\frac{1}{4}+\epsilon}\right)\right)} + \int_0^t Ae^{(t-s)A}W(s) ds$$

$$(-A)^{\frac{1}{4}+\frac{\epsilon}{2}} \int_0^t Ae^{(t-s)A}W(s) ds = \int_0^t \underbrace{(-A)^{1-\frac{\epsilon}{2}} e^{(t-s)A}}_{\|\cdot\|_{\mathcal{L}(H,H)} \leq C/(t-s)^{1-\epsilon/2}} \underbrace{(-A)^{\frac{1}{4}+\epsilon} W(s)}_{\in L^\infty(0, T; H)} ds.$$

Auxiliary Navier-Stokes type equations

$$v(t) = u(t) - z(t), \quad \pi(t) = p(t) - q(t)$$

satisfy

$$\begin{aligned} \partial_t v + (v + z) \cdot \nabla (v + z) + \nabla \pi &= \nu \Delta v + f \\ \operatorname{div} v &= 0 \end{aligned}$$

$$v|_{\partial D} = 0$$

$$v|_{t=0} = v_0.$$

Definition

Given $v_0 \in H$, $f \in L^2(0, T; V')$ and $z \in L^4(0, T; \mathbb{L}^4)$, we say that

$$v \in C([0, T]; H) \cap L^2(0, T; V)$$

is a weak solution of the modified NS equations if

$$\begin{aligned} & \langle v(t), \phi \rangle - \int_0^t b(v(s) + z(s), \phi, v(s) + z(s)) ds \\ &= \langle v_0, \phi \rangle + \int_0^t \langle u(s), A\phi \rangle ds + \int_0^t \langle f(s), \phi \rangle ds \end{aligned}$$

for every $\phi \in D(A)$.

Theorem

For every $v_0 \in H$, $f \in L^2(0, T; V')$ and $z \in L^4(0, T; \mathbb{L}^4)$, there exists a unique weak solution of the modified NS equations. It satisfies

$$\begin{aligned} & \|v(t)\|_{L^2}^2 + 2\nu \int_0^t \|\nabla v(s)\|_{L^2}^2 ds \\ = & \|v_0\|_{L^2}^2 + 2 \int_0^t \langle f(s), v(s) \rangle ds \\ & + 2 \int_0^t (b(v, v, z) + b(z, v, v) + b(z, v, z))(s) ds. \end{aligned}$$

Step 1 (uniqueness). Let $v^{(i)}$ be two solutions. The function $w = v^{(1)} - v^{(2)}$ satisfies

$$\begin{aligned} & \langle w(t), \phi \rangle - \int_0^t \left(b(v^{(1)} + z, \phi, v^{(1)} + z) - b(v^{(2)} + z, \phi, v^{(2)} + z) \right) ds \\ &= \int_0^t \langle w(s), A\phi \rangle ds \end{aligned}$$

for every $\phi \in D(A)$. A simple manipulation gives us

$$\begin{aligned} & b(v^{(1)} + z, \phi, v^{(1)} + z) - b(v^{(2)} + z, \phi, v^{(2)} + z) - b(w, \phi, w) \\ &= b(v^{(2)} + z, \phi, w) + b(w, \phi, v^{(2)} + z) \end{aligned}$$

hence

$$\begin{aligned} \langle w(t), \phi \rangle - \int_0^t b(w, \phi, w) ds &= \int_0^t \langle w, A\phi \rangle ds + \int_0^t \langle \tilde{f}, \phi \rangle ds \\ \tilde{f} &= -B(v^{(2)} + z, w) - B(w, v^{(2)} + z). \end{aligned}$$

By the deterministic Theorem,

$$\begin{aligned} & \|w(t)\|_{L^2}^2 + 2\nu \int_0^t \|\nabla w(s)\|_{L^2}^2 ds \\ &= -2 \int_0^t \left(b(v^{(2)} + z, w, w) - b(w, w, v^{(2)} + z) \right) ds. \end{aligned}$$

From a technical Lemma:

$$\begin{aligned} \left| b(v^{(2)} + z, w, w) \right| &\leq \left| b(v^{(2)}, w, w) \right| + |b(z, w, w)| \\ &\leq 4\epsilon \|w\|_V^2 + \frac{C^2}{\epsilon^3} \|w\|_H^2 \left(\|v^{(2)}\|_{L^4}^4 + \|z\|_{L^4}^4 \right). \end{aligned}$$

Choose $4\epsilon = \nu$.

We get

$$\begin{aligned} & \|w(t)\|_{L^2}^2 + \nu \int_0^t \|\nabla w(s)\|_{L^2}^2 ds \\ &= C \int_0^t \|w(s)\|_H^2 \left(1 + \|v^{(2)}(s)\|_{L^4}^4 + \|z(s)\|_{L^4}^4 \right) ds. \end{aligned}$$

We conclude $w = 0$ by Gronwall lemma, using the assumption

$$z \in L^4(0, T; \mathbb{L}^4)$$

and the property

$$v^{(2)} \in L^4(0, T; \mathbb{L}^4)$$

coming from the regularity

$$v^{(2)} \in C([0, T]; H) \cap L^2(0, T; V).$$

Step 2 (existence). Define the sequence (v^n) by setting $v^0 = 0$ and for every $n \geq 0$, given $v^n \in C([0, T]; H) \cap L^2(0, T; V)$, let v^{n+1} be the solution of the classical NS equations with initial condition v_0 and with

$$f + B(v^n, z) + B(z, v^n) + B(z, z)$$

in place of f . In particular

$$\begin{aligned} & \langle v^{n+1}(t), \phi \rangle - \int_0^t b(v^{n+1}(s), \phi, v^{n+1}(s)) ds \\ &= \langle v_0, \phi \rangle + \int_0^t \langle v^{n+1}(s), A\phi \rangle ds + \int_0^t \langle f(s), \phi \rangle ds \\ & \quad - \int_0^t \langle (B(v^n, z) + B(z, v^n) + B(z, z))(s), \phi \rangle ds \end{aligned}$$

for every $\phi \in D(A)$.

From the deterministic theorem:

$$\begin{aligned} & \|v^{n+1}(t)\|_{L^2}^2 + 2\nu \int_0^t \|\nabla v^{n+1}(s)\|_{L^2}^2 ds \\ = & \|v_0\|_{L^2}^2 + 2 \int_0^t \langle f(s), v^{n+1}(s) \rangle ds \\ & + 2 \int_0^t (b(v^n, v^{n+1}, z) + b(z, v^{n+1}, v^n) + b(z, v^{n+1}, z))(s) ds. \end{aligned}$$

With technical work, for T small enough, we prove iteratively a uniform bound:

$$\sup_{t \in [0, T]} \|v^n(t)\|_H^2 \leq R, \quad \int_0^T \|v^n(s)\|_V^2 ds \leq R$$

Using the previous bounds and similar arguments one can prove that the sequence (v^n) is Cauchy in $C([0, T]; H) \cap L^2(0, T; V)$ (strong topology! Compare with compactness criteria).

The limit v has the right regularity and is a weak solution of the modified NS equations. And satisfies the energy bound.

Remark:

- $W \mapsto z$ measurable
- $(v_0, f, z) \mapsto v$ measurable
- $u = z + v$ implies that $(u_0, f, W) \mapsto u$ measurable.